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## DIFFERENTIAL EQUATIONS FOR INDICATRICES, SPACELIKE AND TIMELIKE CURVES

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**ABSTRACT.** Motivated by the recent work of Deshmukh et al. [20], in this paper we show that Tangent, Binormal, and Principal Normal indicatrices do not form non-trivial differential equations. Finally, we obtain the 4th-order differential equations for spacelike and timelike curves.

*Key words and phrases:* Helix; Slant helix; Tangent, Binormal, Principal Normal indicatrices; Spacelike curves; and Timelike curves.

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## 1. INTRODUCTION

In differential geometry, the curve is among one of the fascinating topics. Helices, spherical curves, and rectifying curves are a few important types of curves that appear in many important applications. For example, helical structures arise in seashells, vines, carbon nanotubes, DNA double, and nano-springs, etc. Though many authors [2, 3, 7, 12, 13, 16, 19, 23, 15] studied curves from the last several decades nevertheless curves are still a relevant and significant area of the research. In the study of curves, the notion of associated curves is pretty exciting. If there exist a mathematical relation between two or more curves, then the curves are known as associated curves.

Izumiya et al. [21] introduced some special curves which are known as a slant helix and conical geodesic curves in Euclidean 3-space. Besides, Izumiya and Takeuchi gave some classifications of the special developable surfaces and obtained an example of a slant helix. In [11], Kula et al. studied the spherical images of the tangent indicatrix and binormal indicatrix of a slant helix. Moreover, they obtained that the spherical images of the slant helices are spherical helices and a curve of constant precession is a slant helix.

In [1], Ali obtained the position vector of a general helix ( $\tau/\kappa = m$ ) associated with Frenet frame and represented the general helix in terms of curvature ( $\kappa$ ) and torsion ( $\tau$ ) through a standard frame of Euclidean 3-space, where  $m$  is a constant given by  $m = \cos[\phi]$ , here  $\phi$  denotes the angle between the axis of a general helix and the tangent of the curve. In [2], Ali et al. extended the concept of a slant helix to Euclidean space of dimension  $n$ , and gave the necessary and sufficient conditions for a curve in Euclidean  $n$ -space to be a slant helix. Moreover, Ali also gave an example of a slant helix in Euclidean 5-space.

Recently, Sahiner [3] defined the associated curves as integral curves of a vector field produced by Frenet vectors of the tangent indicatrix of a curve in Euclidean 3-space and obtained some relations between curvatures and Frenet vectors. Besides, he gave a few techniques to obtain helices and slant helices from special spherical curves and constructed some examples of it. In [4], B. Y. Chen investigated the characterization and classification of the rectifying curves. On the other hand in [5], B. Y. Chen studied via rectifying curves that all geodesics on an arbitrary cone in Euclidean space of dimension 3, are not necessarily a circular cone.

In [6], Yilmaz et al. used the system of linear ordinary differential equations to construct the slant helices. Also, using integration in Minkowski 3-space, they obtained the position vectors for slant helices. In [7], Camci et al. studied and obtained a spherical slant helix and gave some examples of the spherical slant helices in Euclidean 3-space. In [8], Arroyo et al. investigated the unit speed curves contained in a real space form of arbitrary dimension  $m$ . Moreover, they gave a classification of semi-Riemannian Hopf cylinders of  $H_1^3(-1)$  and Hopf cylinders of  $S^3$  with proper mean curvature function.

In [9], Choi et al. introduced the concept of the principal-direction curve and principal-donor curve of a Frenet curve in Euclidean 3-space. Moreover, Choi et al. constructed a canonical method for associated curves and characterized some associated curves in Euclidean 3-space. Kula et al. [10] obtained a relationship between a slant helix and a general helix. Furthermore, Kula et al. deduced some differential equations by characterizing of a slant helix and gave a few examples of slant helices in Euclidean 3-space.

In [17], Lucas et al. studied a weaker version of the classic slant helices in Minkowski 3-space and Euclidean 3-space which are known as general slant helices. Furthermore, Lucas showed that the classic slant helix is a general helix but the converse is not true. Also, he obtained equations that involve the torsion and curvature.

In [19], Deshmukh et al. investigated the rectifying curves via the dilation of the unit speed curve on  $S^2$  (unit sphere) in Euclidean 3-space and obtained a necessary and sufficient condition

for centrode of a unit speed curve in Euclidean 3-space. Moreover, Deshmukh and Chen proved that if a unit speed curve is neither a helix nor a planar curve, then its dilated centrode is always a rectifying curve. Deshmukh et al. [20] shown that for every Frenet curve in Euclidean 3-space, the distance function satisfies a 4th-order differential equation and using this they derived a new characterization of helices. In [22], Ozdemir et al. introduced the notion of type-3 slant helix according to the parallel transport frame in Euclidean 4-space.

Motivated by Deshmukh et al. [20] in this paper, we investigate the distance function. We show that Tangent, Binormal, and Principal Normal indicatrices do not form non-trivial differential equations, and obtain the 4th-order differential equations for spacelike and timelike curves.

## 2. PRELIMINARIES

In this section, we recall some basic concepts of the curves and indicatrices in the Euclidean 3-space. Let  $\beta : I \rightarrow \mathbb{R}^3$  represents the unit speed curve in the Euclidean 3-space and  $T', N', B'$  be the three orthonormal vectors of the Frenet frame  $\{T, N, B\}$ , given by

$$T = \frac{d\beta}{ds}, \quad N = \frac{T'}{\kappa}, \quad B = T \times N$$

where  $T, N, B$  represent the unit tangent vector field, unit principal normal vector field and unit binormal vector field, respectively.

The Serret-Frenet formulae are given by

$$(2.1) \quad \begin{cases} T'(s) = \kappa(s) N(s) \\ N'(s) = -\kappa(s) T(s) + \tau(s) B(s) \\ B'(s) = -\tau(s) N(s) \end{cases}$$

where  $\kappa(s) = \|T'(s)\|$  denote the curvature and  $\tau(s) = -\langle B'(s), N(s) \rangle$  denote the torsion of the curve  $\beta$ . Here the curve  $\beta$  is parameterized in terms of the arc-length parameter  $s$  [18].

If the position vector of the curve  $\beta$  lies in the rectifying plane then the curve is known as a rectifying curve. The distance function  $d(s) = \|\beta(s)\|$  of a rectifying curve  $\beta$  satisfies the following equation

$$d(s) = \sqrt{s^2 + c_1 s + c_2}$$

here  $c_1$  and  $c_2$  denote the arbitrary constants.

Furthermore, it can be shown that the unit speed curve  $\beta$  is also a rectifying curve if and only if the ratio of torsion  $\tau$  and curvature  $\kappa$  verifies

$$\frac{\tau}{\kappa} = as + b$$

where  $a \neq 0$  and  $b$  are constants [4].

Choi and Kim investigated the relationship between curvature and torsion of the principal-direction curve and principal-donor curve in [9].

**Theorem 2.1.** [9] *Let  $\beta$  be a Frenet curve in Euclidean 3-space with the curvature  $\kappa$  and the torsion  $\tau$  and  $\bar{\beta}$  be the principal-direction curve of the curve  $\beta$ . Then the curvature  $\bar{\kappa}$  and torsion  $\bar{\tau}$  of the principal-direction curve  $\bar{\beta}$  are given by*

$$\bar{\kappa} = \sqrt{\kappa^2 + \tau^2} \quad \text{and} \quad \bar{\tau} = \frac{\kappa^2}{\kappa^2 + \tau^2} \left( \frac{\tau}{\kappa} \right)'$$

**Theorem 2.2.** [9] *Let  $\beta$  be a principal-donor curve of the curve  $\bar{\beta}$  in Euclidean 3-space with the curvature  $\bar{\kappa}$  and torsion  $\bar{\tau}$ . Then the curvature  $\kappa$  and torsion  $\tau$  of the principal-donor curve  $\beta$  are given by*

$$\kappa = \bar{\kappa} \left| \cos \left( \int \bar{\tau} ds \right) \right| \quad \text{and} \quad \tau = \bar{\kappa} \sin \left( \int \bar{\tau} ds \right)$$

A curve  $\beta$  is said to be general helix if unit tangent  $T(s)$  makes a constant angle with a fixed straight line. Likewise, if unit principal normal  $N(s)$  makes a constant angle with a fixed straight line then a curve  $\beta$  is said to be slant helix.

Let  $\beta$  be a unit speed curve in Euclidean space with Frenet vectors  $T$ ,  $N$  and  $B$ . The unit tangent vectors along the curve  $\beta$  generate a curve  $\beta_t$  on the unit sphere centered at the origin, called the tangent indicatrix of curve  $\beta$ . Similarly, we have the binormal indicatrix  $\beta_b$  and principal normal indicatrix  $\beta_n$  [10].

Deshmukh and B. Y. Chen shown that for every Frenet curve in Euclidean 3-space, the distance function satisfies a general differential equation. We recall the following proposition from [20].

**Proposition 2.3.** *If  $\beta$  be a unit speed curve then every unit speed Frenet curve satisfies the following equation:*

$$(2.2) \quad \rho\sigma h''' + (\rho\sigma' + 2\rho'\sigma)h'' + \left\{ (\sigma\rho')' + \frac{\rho}{\sigma} + \frac{\sigma}{\rho} \right\} h' + \left( \frac{\sigma}{\rho} \right)' h = (\sigma\rho')' + \frac{\rho}{\sigma},$$

where  $\rho = \kappa^{-1}$ ,  $\sigma = \tau^{-1}$ ,  $h(s) = d(s)d'(s)$ .

The Minkowski 3-space  $\mathbb{E}_1^3$  is the Euclidean 3-space provided with the standard flat metric given by

$$g = -dx_1^2 + dx_2^2 + dx_3^2,$$

where  $(x_1, x_2, x_3)$  is a rectangular coordinate system of  $\mathbb{E}_1^3$ .

Since  $g$  is an indefinite metric, recall that a vector  $v \in \mathbb{E}_1^3$  can have one of the three causal characters; it can be spacelike if  $g(v, v) > 0$  or  $v = 0$ , timelike if  $g(v, v) < 0$  and lightlike (null) if  $g(v, v) = 0$  and  $v \neq 0$ . Analogously, an arbitrary curve  $\beta = \beta(s)$  in  $\mathbb{E}_1^3$  can locally be spacelike, timelike or lightlike (null), if all of its velocity vectors  $\beta'(s)$  are respectively spacelike, timelike or lightlike. The norm of a vector  $v$  is given by  $\|v\| = \sqrt{|g(v, v)|}$  and the spacelike or (timelike) curve  $\beta(s)$  is said to be of unit speed if  $g(\beta'(s), \beta'(s)) = \pm 1$  [14].

### 3. DERIVATION OF THE DIFFERENTIAL EQUATIONS

In this section, first we give some propositions for indicatrices, Serret-Frenet formulae, and a few useful results for spacelike and timelike curves. Finally, we obtain the 4th-order differential equations for spacelike and timelike curves.

**Proposition 3.1.** *If  $\beta$  be a unit speed curve then tangent indicatrix  $\beta_t$  of the curve  $\beta$  does not form a non-trivial differential equation.*

*Proof.* Since the tangent indicatrix  $\beta_t$  has constant norm equal to one. By differentiating the distance function  $d(s) = \|\beta_t(s)\|$ , we get  $d'(s) = 0$ .

**Proposition 3.2.** *If  $\beta$  be a unit speed curve then binormal indicatrix  $\beta_b$  of curve  $\beta$  does not form a non-trivial differential equation.*

*Proof.* Since the binormal indicatrix  $\beta_b$  has constant norm equal to one. By differentiating the distance function  $d(s) = \|\beta_b(s)\|$ , we get  $d'(s) = 0$ .

**Proposition 3.3.** *If  $\beta$  be a unit speed curve then principal normal indicatrix  $\beta_n$  of curve  $\beta$  does not form a non-trivial differential equation.*

*Proof.* Since the principal normal indicatrix  $\beta_n$  has constant norm equal to one. By differentiating the distance function  $d(s) = \|\beta_n(s)\|$ , we get  $d'(s) = 0$ .

**Remark 3.1.** Suppose  $\beta$  denote a spacelike curve with a spacelike principal normal  $N$  and  $\beta'$  be the tangent vector field, then the Serret-Frenet formulae are given by

$$(3.1) \quad \begin{cases} T' = \kappa N \\ N' = -\kappa T + \tau B \\ B' = \tau N \end{cases}$$

where  $\langle T, T \rangle = 1$ ,  $\langle N, N \rangle = 1$ ,  $\langle B, B \rangle = -1$ ,  $\langle T, N \rangle = \langle T, B \rangle = \langle N, B \rangle = 0$ .

From the above formula, we have the following

$$(3.2) \quad \begin{cases} \langle \beta, T \rangle' = 1 + \kappa \langle \beta, N \rangle \\ \langle \beta, N \rangle' = -\kappa \langle \beta, T \rangle + \tau \langle \beta, B \rangle \\ \langle \beta, B \rangle' = \tau \langle \beta, N \rangle \end{cases}$$

**Theorem 3.4.** *Suppose  $\beta$  denote a spacelike curve with a spacelike principal normal  $N$ , then the function  $f(s) = d(s)d'(s)$  satisfies the following differential equation*

$$(3.3) \quad \begin{aligned} \frac{f'''}{\tau\kappa} + \left[ \frac{1}{\tau'\kappa} + \frac{2}{\tau\kappa'} \right] f'' + \left[ \frac{1}{\tau'\kappa'} + \frac{1}{\tau\kappa''} + \frac{\kappa}{\tau} - \frac{\tau}{\kappa} \right] f' + \left[ \frac{\kappa}{\tau'} + \frac{\kappa'}{\tau} \right] f \\ = \left[ \frac{1}{\tau\kappa'} \right]' - \frac{\tau}{\kappa} \end{aligned}$$

where  $d(s) = \|\beta(s)\|$  is the distance function of  $\beta$ .

*Proof.* Differentiating  $d^2(s) = \langle \beta(s), \beta(s) \rangle$  and making use of equation (3.1), we get

$$(3.4) \quad f = \langle \beta, T \rangle$$

Now, differentiating above equation and using (3.2), we get

$$(3.5) \quad f' - 1 = \kappa \langle \beta, N \rangle$$

Further, differentiating equation (3.5), yields

$$(3.6) \quad \frac{1}{\tau\kappa} f'' + \frac{1}{\tau\kappa'} f' + \frac{\kappa}{\tau} f - \frac{1}{\tau\kappa'} = \langle \beta, B \rangle$$

Now, differentiating equation (3.6) and using (3.2), (3.5), we get the desired result.

**Remark 3.2.** Suppose  $\beta$  denote a spacelike curve with a timelike principal normal  $N$  and  $\beta'$  be the tangent vector field, then the Serret-Frenet formulae are given by

$$(3.7) \quad \begin{cases} T' = \kappa N \\ N' = \kappa T + \tau B \\ B' = \tau N \end{cases}$$

where  $\langle T, T \rangle = 1$ ,  $\langle N, N \rangle = -1$ ,  $\langle B, B \rangle = 1$ ,  $\langle T, N \rangle = \langle T, B \rangle = \langle N, B \rangle = 0$ .

From the above equation, we get the following

$$(3.8) \quad \begin{cases} \langle \beta, T \rangle' = 1 + \kappa \langle \beta, N \rangle \\ \langle \beta, N \rangle' = \kappa \langle \beta, T \rangle + \tau \langle \beta, B \rangle \\ \langle \beta, B \rangle' = \tau \langle \beta, N \rangle \end{cases}$$

**Theorem 3.5.** Suppose  $\beta$  denote a spacelike curve with a timelike principal normal  $N$ , then the function  $f(s) = d(s)d'(s)$  satisfies the following differential equation

$$(3.9) \quad \frac{f'''}{\tau\kappa} + \left[ \frac{1}{\tau'\kappa} + \frac{2}{\tau\kappa'} \right] f'' + \left[ \frac{1}{\tau'\kappa'} + \frac{1}{\tau\kappa''} - \frac{\kappa}{\tau} - \frac{\tau}{\kappa} \right] f' - \left[ \frac{\kappa}{\tau'} + \frac{\kappa'}{\tau} \right] f = \left[ \frac{1}{\tau\kappa'} \right]' - \frac{\tau}{\kappa}$$

where  $d(s) = \|\beta(s)\|$  is the distance function of  $\beta$ .

*Proof.* Differentiating  $d^2(s) = \langle \beta(s), \beta(s) \rangle$  and making use of equation (3.7), we get

$$(3.10) \quad f = \langle \beta, T \rangle$$

Using (3.10) and (3.8), a simple computation gives

$$(3.11) \quad f' - 1 = \kappa \langle \beta, N \rangle$$

Now, differentiating (3.11), we get

$$(3.12) \quad \frac{1}{\tau\kappa} f'' + \frac{1}{\tau\kappa'} f' - \frac{\kappa}{\tau} f - \frac{1}{\tau\kappa'} = \langle \beta, B \rangle$$

Finally, differentiating equation (3.12) and using (3.8), (3.11), we get the desired result.

**Remark 3.3.** Suppose  $\beta$  denote a spacelike curve with a lightlike principal normal  $N$  and  $\beta'$  be the tangent vector field, then the Serret-Frenet formulae are given by

$$(3.13) \quad \begin{cases} T' = \kappa N \\ N' = \tau N \\ B' = -\kappa T - \tau B \end{cases}$$

where  $\langle T, T \rangle = 1, \langle N, B \rangle = 1, \langle N, N \rangle = \langle B, B \rangle = \langle T, N \rangle = \langle T, B \rangle = 0$ .

From the above formula, we have the following

$$(3.14) \quad \begin{cases} \langle \beta, T \rangle' = 1 + \kappa \langle \beta, N \rangle \\ \langle \beta, N \rangle' = \tau \langle \beta, N \rangle \\ \langle \beta, B \rangle' = -\kappa \langle \beta, T \rangle - \tau \langle \beta, B \rangle \end{cases}$$

**Theorem 3.6.** Suppose  $\beta$  denote a spacelike curve with a lightlike principal normal  $N$ , then the function  $f(s)$  satisfies the following differential equation

$$(3.15) \quad \frac{f'''}{\tau\kappa} + \left[ \frac{1}{\tau'\kappa} + \frac{2}{\tau\kappa'} - \frac{1}{\kappa} \right] f'' + \left[ \frac{1}{\tau'\kappa'} + \frac{1}{\tau\kappa''} - \frac{1}{\kappa'} \right] f' = \left[ \frac{1}{\tau\kappa'} \right]' - \frac{1}{\kappa'}$$

where  $f(s) = d(s)d'(s)$ , and  $d(s) = \|\beta(s)\|$  is the distance function of the curve  $\beta$ .

*Proof.* Differentiating  $d(s) = \|\beta(s)\|$  and using equation (3.13), we get

$$(3.16) \quad f = \langle \beta, T \rangle$$

From equations (3.16) and (3.14), we have

$$(3.17) \quad f' - 1 = \kappa \langle \beta, N \rangle$$

Now, differentiating (3.17), yields

$$(3.18) \quad \frac{1}{\tau\kappa} f'' + \frac{1}{\tau\kappa'} f' - \frac{1}{\tau\kappa'} = \langle \beta, N \rangle$$

Differentiating equation (3.18) and using (3.16), we get the result.

**Remark 3.4.** Suppose  $\beta$  denote a timelike curve and  $\beta'$  be the tangent vector field, then the Serret-Frenet formulae are given by

$$(3.19) \quad \begin{cases} T' = \kappa N \\ N' = \kappa T + \tau B \\ B' = -\tau N \end{cases}$$

where  $\langle T, T \rangle = -1$ ,  $\langle N, N \rangle = 1$ ,  $\langle B, B \rangle = 1$ ,  $\langle T, N \rangle = \langle T, B \rangle = \langle N, B \rangle = 0$ .

From the above formula, we have the following

$$(3.20) \quad \begin{cases} \langle \beta, T \rangle' = -1 + \kappa \langle \beta, N \rangle \\ \langle \beta, N \rangle' = \kappa \langle \beta, T \rangle + \tau \langle \beta, B \rangle \\ \langle \beta, B \rangle' = -\tau \langle \beta, N \rangle \end{cases}$$

**Theorem 3.7.** Suppose  $\beta$  denote a timelike curve, then the function  $f(s) = d(s)d'(s)$  satisfies the following differential equation

$$(3.21) \quad \begin{aligned} \frac{f'''}{\tau\kappa} + \left[ \frac{1}{\tau'\kappa} + \frac{2}{\tau\kappa'} \right] f'' + \left[ \frac{1}{\tau'\kappa'} + \frac{1}{\tau\kappa''} - \frac{\kappa}{\tau} + \frac{\tau}{\kappa} \right] f' - \left[ \frac{\kappa}{\tau'} + \frac{\kappa'}{\tau} \right] f \\ = - \left[ \frac{1}{\tau\kappa'} \right]' - \frac{\tau}{\kappa} \end{aligned}$$

where  $d(s) = \|\beta(s)\|$  is the distance function of  $\beta$ .

*Proof.* Differentiating  $d^2(s) = \langle \beta(s), \beta(s) \rangle$  and making use of equation (3.19), we get

$$(3.22) \quad f = \langle \beta, T \rangle$$

Now, differentiating above equation and using (3.20), we get

$$(3.23) \quad f' + 1 = \kappa \langle \beta, N \rangle$$

Further, differentiating equation (3.23), yields

$$(3.24) \quad \frac{1}{\tau\kappa} f'' + \frac{1}{\tau\kappa'} f' - \frac{\kappa}{\tau} f + \frac{1}{\tau\kappa'} = \langle \beta, B \rangle$$

Now, differentiating equation (3.24) and using (3.20), (3.23), the result follows.

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